



# High permittivity and low dielectric loss of the $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$ ceramics

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## ABSTRACT

The  $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  (CSCTO) giant dielectric ceramics were prepared by conventional solid-state method. X-ray diffraction patterns revealed that a small amount of  $\text{Sr}^{2+}$  ( $x < 0.2$ ) had no obvious effect on the phase structure of the CSCTO ceramics, while with increasing the  $\text{Sr}^{2+}$  content, a second phase of  $\text{SrTiO}_3$  appeared. Electrical properties of CSCTO ceramics greatly depended on the  $\text{Sr}^{2+}$  content. The  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics exhibited a higher permittivity (71,153) and lower dielectric loss (0.022) when measured at 1 kHz at room temperature. The ceramics also performed good temperature stability in the temperature range from  $-50^\circ\text{C}$  to  $100^\circ\text{C}$  at 1 kHz. By impedance spectroscopy analysis, all compounds were found to be electrically heterogeneous, showing semiconducting grains and insulating grain boundaries. The grain resistance was  $1.28\ \Omega$  and the grain boundary resistance was  $1.31 \times 10^5\ \Omega$ . All the results indicated that the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics were very promising materials with higher permittivity for practical applications.

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## 1. Introduction

Perovskite  $\text{CaCu}_3\text{Ti}_4\text{O}_{12}$  (CCTO) has attracted much attention because of its intriguing mechanism and potential technical applications in microelectronic devices. Since  $\text{CaCu}_3\text{Ti}_4\text{O}_{12}$  (CCTO) giant dielectric material was reported by Subramanian et al. in 2000, CCTO has attracted considerable attention owing to its high dielectric permittivity ( $\sim 10^4$ – $10^5$ ) at low frequencies and room temperature [1–6]. CCTO simultaneously has good temperature stability over a wide temperature ranging from 100 K to 600 K [7]. At temperature below 100 K, the dielectric constant drops rapidly to around 100 without any structural phase transition [8,9]. Despite the fact that much research work has been done, including the experiment and theory, to reveal the origin of the giant dielectric permittivity of CCTO ceramics, it still remains controversial and unsolved problems to date. For CCTO ceramics, the internal barrier layer capacitance (IBLC) model is commonly accepted as the nature of its giant dielectric response. Sinclair et al. reported that CCTO ceramics consisting of semiconducting grains and insulating grain boundaries were electrically heterogeneous [8].

Up to now, several investigation on the CCTO ceramics and single crystal as well as other related materials have been reported, but the high loss tangent ( $\tan \delta$ ) of the CCTO ceramics ( $\tan \delta > 0.05$  at 1 kHz) is still an obstruction to the technical use of the CCTO. Number of attempts has been made to decrease  $\tan \delta$  by various

compositional modifications, such as  $\text{TeO}_2$ -doped CCTO [10],  $\text{ZrO}_2$ -doped CCTO [11], CCTO/ $\text{CaTiO}_3$  composite [12,13],  $\text{TiO}_2$ -rich CCTO [14], and the substitution of La for Cu in CCTO [15] have also been reported to be effective in reducing the dielectric loss and still maintaining the high dielectric constant.

In this work, the cations substitution in the A site of CCTO by  $\text{Sr}^{2+}$  has been studied. The phase structure, microstructure, electrical properties of the CSCTO ceramics were investigated.

## 2. Experimental

$\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ( $x = 0, 0.10, 0.15, 0.20, 0.25$ ) ceramics were prepared by the conventional solid-state method.  $\text{CaCO}_3$  (99.99%),  $\text{SrCO}_3$  (99.99%),  $\text{CuO}$  (99.99%), and  $\text{TiO}_2$  (99.99%) were used as starting materials. The stoichiometric powders were mixed by ball-milled in ethanol for 10 h using agate balls. The mixed powders were calcined at  $900^\circ\text{C}$  for 10 h in air and then pressed into disks with a diameter of 15 mm under 100 MPa pressure using a solution of poly vinyl alcohol (PVA) as binder. The pellets were sintered at  $1060^\circ\text{C}$ ,  $1080^\circ\text{C}$ ,  $1090^\circ\text{C}$  and  $1100^\circ\text{C}$  for 5 h in air, respectively.

The crystalline phase of the sintered specimens was identified by X-ray diffraction (XRD, D/max-2550/PC, Rigaku, Japan) technique using  $\text{Cu K}\alpha$  radiation, and the surface microstructure was observed by scanning electron microscopy (SEM, Quanta 200, Philips, Netherlands). In order to measure the electric properties, silver paste was painted on both sides of samples as the electrodes and were fired at  $840^\circ\text{C}$ , and then the samples were measured by Agilent 4294A impedance analyzer in the frequency ranging from 40 Hz to 110 MHz.

## 3. Results and discussion

Fig. 1 shows the XRD patterns of CSCTO powders with different  $\text{Sr}^{2+}$  contents calcined at  $900^\circ\text{C}$  for 10 h. It can be seen from Fig. 1 that the samples of  $x < 0.2$  are single CCTO phase, which could be

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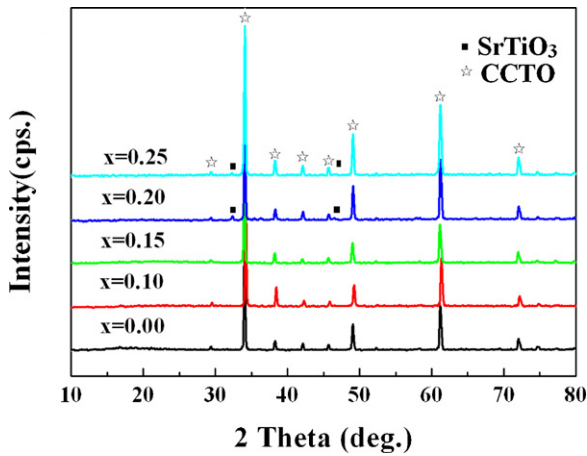


Fig. 1. X-ray diffraction patterns of  $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  as a function of  $\text{Sr}^{2+}$  doping.

indexed to a cubic perovskite-related structure according to JCPDS-75-2188 [16], indicating that a little amount of  $\text{Sr}^{2+}$  can enter into the lattice of CCTO. For the samples of  $x \geq 0.2$ , some diffraction peaks of the second phase were found around  $2\theta \sim 32.5^\circ, 46.7^\circ$ , which means a small amount of  $\text{SrTiO}_3$  impurity phase starts to appear. The ionic radius of  $\text{Sr}^{2+}$  ( $r_{\text{Sr}^{2+}} = 1.44 \text{ \AA}$ ) is close to that of  $\text{Ca}^{2+}$  ( $r_{\text{Ca}^{2+}} = 1.34 \text{ \AA}$ ), then small amount of  $\text{Sr}^{2+}$  can easily enter into the

structure lattice and form a solid solution, while with increasing the  $\text{Sr}^{2+}$  content, the second phase of  $\text{SrTiO}_3$  appears.

Fig. 2 demonstrates the SEM images of CSCTO ceramics sintered at  $1080^\circ\text{C}$  with various  $x$  values. Abnormal grain growth with grain size larger than  $100 \mu\text{m}$  can be found in Fig. 2(a)–(c). With increasing  $\text{Sr}^{2+}$  content, the average grain size of CSCTO ceramics increases firstly and then decreases. The average grain size of  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics is about  $150\text{--}200 \mu\text{m}$ , which is greater than that of other doped samples. In addition, normal grain growth, clear grain boundaries and the porosities can be obtained only in Fig. 2(d), which shows that further increasing  $x$  to 0.25 can lead to the decrease of grain size and disappearance of abnormal grain growth.

Fig. 3(a) illustrates the frequency dependence of the permittivity ( $\epsilon_r$ ) of the CSCTO ceramics with different  $x$  values measured at room temperature. It exhibits the plateau permittivity of the CSCTO ceramics at the frequency ranging from 40 Hz to 100 kHz. When the frequency surpasses 100 kHz, the permittivity drastically decreases and then almost turns to be a small constant. The results show that the CSCTO ceramics have good frequency stability. This is a typical characteristic of Debye relaxation, and it is consistent with the reported results [17–20]. In addition, the permittivity of ceramics with  $x = 0.15\text{--}0.25$  (below 13,000) is smaller than the values of sample with  $x = 0, 0.1$  (above 70,000). It can also be seen that the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics have higher permittivity (71,153) at 1 kHz than that of the CCTO (61,425) ceramics. The giant permittivity can be explained as an internal barrier layer capacitance (IBLC)

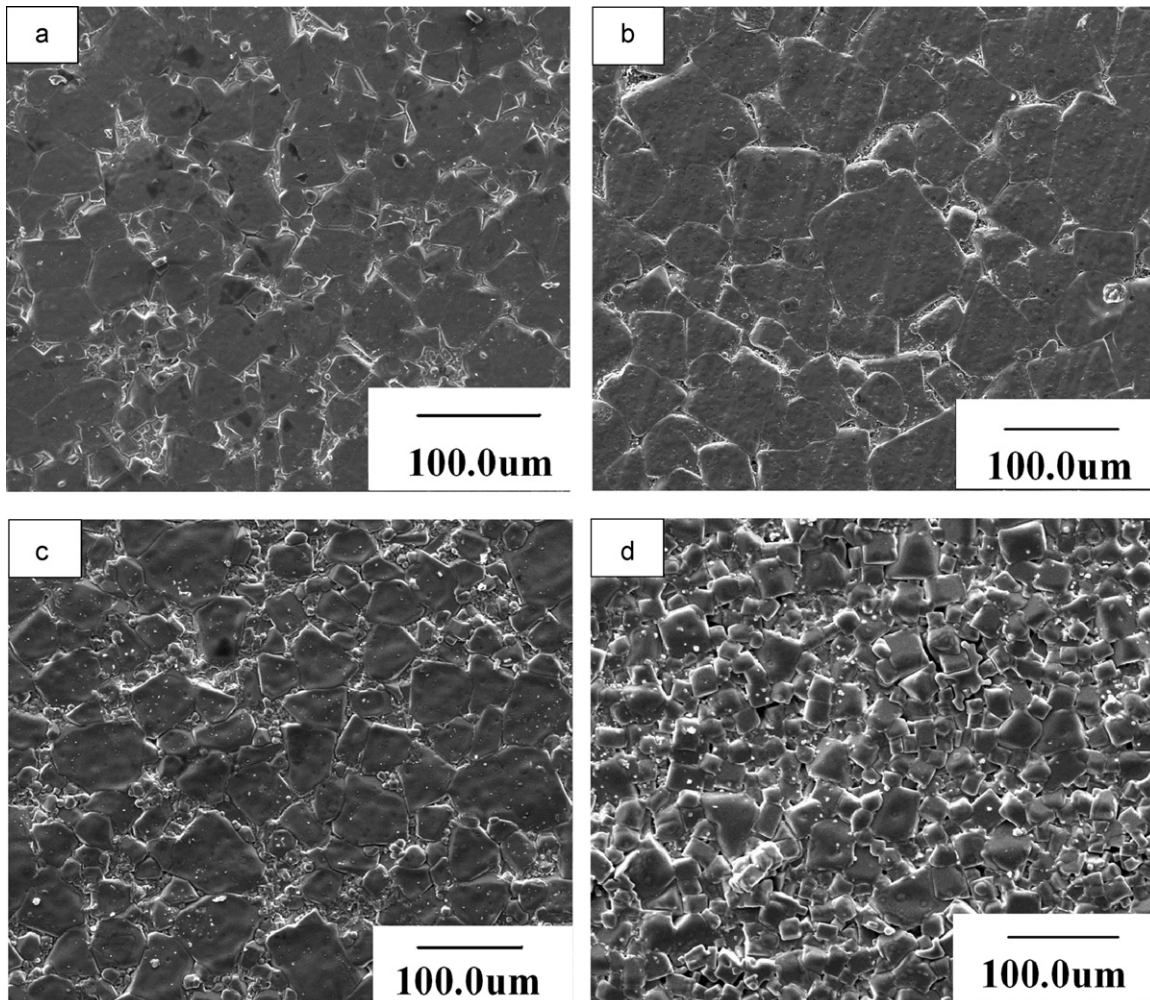
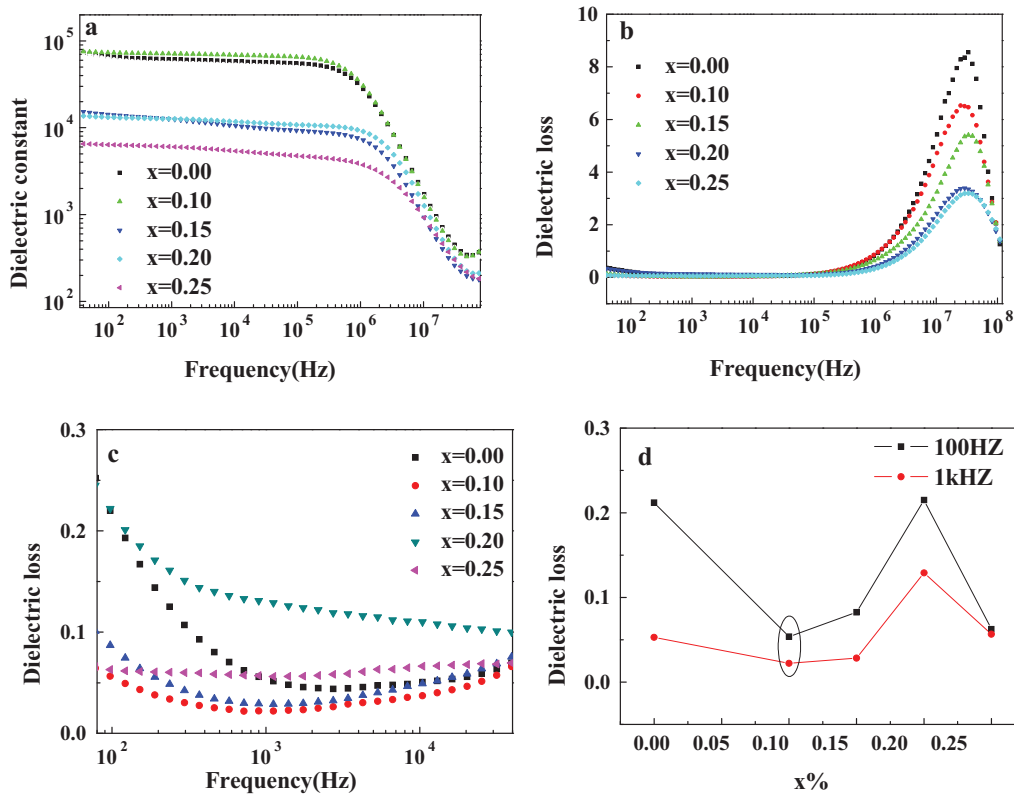


Fig. 2. Surface SEM images of  $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics: (a)  $x = 0.00$ , (b)  $x = 0.10$ , (c)  $x = 0.15$ , (d)  $x = 0.25$ .



**Fig. 3.** Frequency dependence of permittivity (a) and dielectric loss (b) of  $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics at room temperature, Fig. 3(c) is a magnified view of dielectric loss from 100 Hz to 40 kHz, the dielectric loss at 100 Hz and 1000 Hz with different  $\text{Sr}^{2+}$  content is shown in Fig. 3(d).

effect by Ferrarelli et al. [21]. Consequently, appropriately substituting  $\text{Ca}^{2+}$  with  $\text{Sr}^{2+}$  in CCTO can improve the dielectric properties of the ceramics.

The frequency dependence of dielectric loss of the CSCTO ceramics as a function of  $x$  measured at room temperature is shown in Fig. 3(b). It is evident that the dielectric loss has good frequency stability from 40 Hz to 100 kHz. As the frequency increases above 100 kHz, dielectric loss drastically increases and then decreases. It can also be found that the dielectric loss of the CSCTO ceramics gradually reduces with increasing the  $\text{Sr}^{2+}$  content in the frequency from 1.6 MHz to 56 MHz. Fig. 3(c) is a magnified view of dielectric loss from 100 Hz to 40 kHz, dielectric loss decreases and then increases with the  $\text{Sr}^{2+}$  content in Fig. 3(c). The dielectric loss at 100 Hz and 1 kHz with different  $\text{Sr}^{2+}$  content is shown in Fig. 3(d). It can be seen from Fig. 3(d) that the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics possess the lowest dielectric loss (0.0537) at 100 Hz and (0.022) at 1 kHz. With reference to Fig. 3(a) and (b), the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics have higher permittivity (about 71,153) and lower dielectric loss (about 0.022) at 1 kHz at room temperature.

It is well known that the impedance spectroscopy (IS) is a powerful tool in separating out the grain and the grain boundary effects. For CCTO ceramics, an impedance spectroscopy analysis demonstrates that CCTO ceramics are electrically heterogeneous, consisting of semiconducting grains and insulating grain boundaries, which are simulated by Zsimp Win software. The impedance of electroceramics has been modeled as an equivalent circuit consisting of two parallel RC elements connected in series, of which one RC element corresponds to the semiconducting grains and the other corresponds to the insulating grain boundaries. For such a circuit, each RC element ideally gives rise to semicircular arc in complex impedance plane  $Z^*$  [22], where

$$Z^* = Z' - iZ''$$

where

$$Z' = \frac{R_g}{1 + (\omega R_g C_g)^2} + \frac{R_{gb}}{1 + (\omega R_{gb} C_{gb})^2}$$

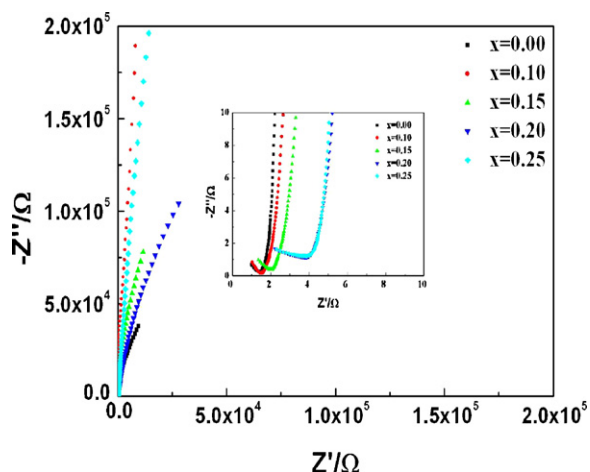
and

$$Z'' = R_g \left[ \frac{\omega R_g C_g}{1 + (\omega R_g C_g)^2} \right] + R_{gb} \left[ \frac{\omega R_{gb} C_{gb}}{1 + (\omega R_{gb} C_{gb})^2} \right]$$

$R_g$  and  $C_g$  represent the grain resistance and capacitance, respectively, while  $R_{gb}$  and  $C_{gb}$  represent the grain boundary resistance and capacitance, respectively.

Fig. 4 shows complex impedance spectroscopy of CSCTO ceramics with various  $x$  values. The  $R_g$  and  $R_{gb}$  values can be obtained by an impedance spectrum analysis. The data of all CSCTO ceramics shows a single semicircular arc with a nonzero high frequency intercept. There are two semicircular arcs in the complex impedance spectroscopy in the measuring frequency ranging from 40 Hz to 110 MHz, corresponding to the two parallel RC elements in series, which indicates that the higher permittivity is due to IBLC model [23,24]. The IS results show the high frequency impedance corresponds to grain resistance  $R_g$  and the low frequency impedance corresponds to grain boundary resistance  $R_{gb}$ . As the frequency increases, grain boundary resistance ( $R_{gb}$ ) and grain resistance ( $R_g$ ) first increase and then decrease. As is shown in Fig. 4, the grain resistance ( $R_g$ ) of the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics is  $1.28 \Omega$  and the grain boundary resistance ( $R_{gb}$ ) is  $1.31 \times 10^5 \Omega$ . Generally, the grain resistance  $R_g$  is much smaller than the grain boundary resistance  $R_{gb}$ . These results indicate that the grains are semiconducting and the grain boundaries are insulating, which is consistent with the internal barrier layer capacitance (IBLC) model and demonstrates that CSCTO ceramics are electrically heterogeneous. The  $R_{gb}$  has the similar change trend as  $R_g$  and the maximum





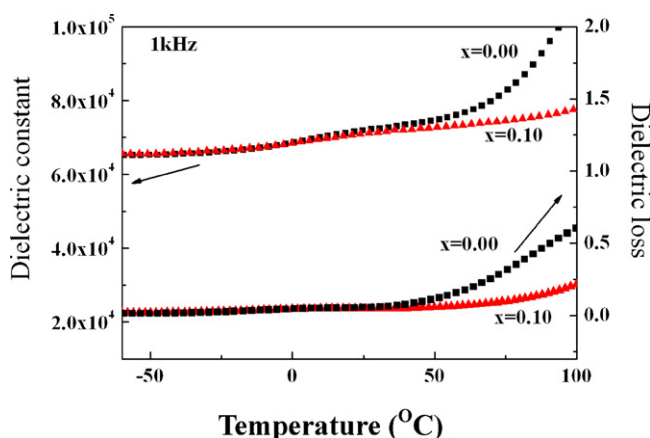
**Fig. 4.** The Cole–Cole plot of impedance at room temperature for  $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics. The inset shows an enlarged view for the highest frequency data close to the origin.

value of  $R_{gb}$  and  $R_g$  of the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics can be seen in Fig. 4. Experimental results are basically consistent with the reduction in dielectric loss, the dielectric loss is approximately written by Eq. (1) [25]

$$\tan \delta = \frac{1}{wR_{gb}C} + wR_gC \quad (1)$$

where  $C$  and  $w$  represent measured capacitance and angle frequency, respectively. It can be seen that the first term on the right of Eq. (1) has higher weight at low frequencies, while the second term becomes predominant at high frequencies, due to  $R_{gb} \geq R_g$ . According to Eq. (1),  $\tan \delta \approx 1/wR_{gb}C$  at low frequencies, while  $w > 10^5$  Hz,  $\tan \delta \approx wR_gC$ . This is in conformity with the decrease of dielectric loss in Fig. 3(b) and (c). Based on IBL model, this phenomenon suggests that the  $\text{Sr}^{2+}$  affects the electrical properties of CCTO.

Fig. 5 shows the temperature dependence of the permittivity and dielectric loss of the CCTO and the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics at 1 kHz from  $-50^\circ\text{C}$  to  $100^\circ\text{C}$ . The results reveal that the CCTO and  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics have higher permittivity (above 60,000) and lower dielectric loss (below 0.05) with good temperature stability in the temperature ranging from  $-50^\circ\text{C}$  to  $50^\circ\text{C}$ . Further increasing the temperature leads to the increase in both permittivity and dielectric loss. Compared with CCTO ceramics, the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics have better temperature stability.



**Fig. 5.** Temperature dependence of dielectric properties of  $\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ( $x=0.00, 0.10$ ) specimens at 1 kHz.

## 4. Conclusions

$\text{Ca}_{1-x}\text{Sr}_x\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ( $0 \leq x \leq 0.25$ ) giant dielectric ceramics were prepared by the conventional solid-state method. The phase structure, microstructure and dielectric properties of CSCTO ceramics were investigated. XRD results indicated that when  $x < 0.2$ ,  $\text{Sr}^{2+}$  could enter into the cubic perovskite structure lattice and form a solid solution, while with increasing the  $\text{Sr}^{2+}$  content, the second phase of  $\text{SrTiO}_3$  appeared. Scanning electron microscopy analysis showed that the grains of CSCTO ( $x=0.10$ ) ceramics were greater than that of CCTO and the grains of CSCTO ( $x > 0.10$ ) ceramics gradually decreased with increasing the  $\text{Sr}^{2+}$  content. The permittivity and the dielectric loss of the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics were 71,153 and 0.022, respectively when measured at 1 kHz at room temperature. The permittivity of the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics was more than 65,000 and corresponding dielectric loss was below 0.05 at 1 kHz in the range of  $-50^\circ\text{C}$  to  $100^\circ\text{C}$ . The results indicated that the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics performed higher permittivity, lower dielectric loss and better temperature stability. The experimental results were consistent with the simulated data obtained from Zsimp Win software, which demonstrated that the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics were electrically heterogeneous. All results indicated that the  $\text{Ca}_{0.9}\text{Sr}_{0.1}\text{Cu}_3\text{Ti}_4\text{O}_{12}$  ceramics were very promising materials with higher permittivity and lower dielectric loss for practical applications.

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